# Special motions for spacelike curve in Minkowski 3-space 

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#### Abstract

Existence of acceleration pole points in special Frenet and Bishop motions for spacelike curve with a spacelike binormal in Minkowski 3 -space $E_{1}^{3}$ are dependence into that, the curve $\alpha$ is not a general helix or planar. The ratio of torsion and curvature is by taking as a constant or non constant in our study. Then we show that, if the ratio of curvatures is constant, then there is not acceleration pole points of motion.


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## 1 Preliminaries

Let $R^{3}$ be the real vector space with its usual vector structure. The Minkowski 3-space is the metric space $E_{1}^{3}=\left(R^{3},\langle,\rangle_{L}\right)$, where the metric $\langle,\rangle_{L}$ is given by

$$
\langle x, y\rangle_{L}=x_{1} y_{1}+x_{2} y_{2}-x_{3} y_{3}: x=\left(x_{1}, x_{2}, x_{3}\right), y=\left(y_{1}, y_{2}, y_{3}\right)
$$

The metric $\langle,\rangle_{L}$ is called the Lorentzian metric $[4,6]$.
A vector $x \in E_{1}^{3}$ is called:
i) Spacelike if $\langle x, x\rangle_{L}>0$ or $x=0$,
ii) Timelike if $\langle x, x\rangle_{L}<0$,
iii) Null (lightlike) if $\langle x, x\rangle_{L}=0$ and $x \neq 0$.

Denote by $\{T, N, B\}$ the moving Frenet frame and the moving Bishop frame along the regular curve $\alpha=\alpha(t)$ that are parameterized by the length- arc parameter $t$ The Frenet trihedron consists of the tangent vector $T$, the principle normal vector $N$ and the binormal vector $B$, and the Bishop trihedron consists of the tangent vector $T$, the $1^{\text {st }}$ principle normal vector $N_{1}$ and $2^{\text {nd }}$ principle normal vector $N_{2}$, which are three mutually orthogonal axes.

If $\alpha$ is a spacelike curve with a spacelike binormal, then this set of orthogonal unit vectors, known as the Frenet- serret frame has the following properties:

$$
\begin{aligned}
& \dot{T}=\kappa N, \dot{N}=\kappa T+\tau B, \dot{B}=\tau N \\
& \langle T, T\rangle_{L}=1,\langle N, N\rangle_{L}=-1,\langle B, B\rangle_{L}=1
\end{aligned}
$$

[1,5,6]. In this formulas, the normal vector is $N=\frac{\dot{T}}{\kappa}$, where $\kappa=\sqrt{-\langle\dot{T}, \dot{T}\rangle_{L}}$ is the curvature of $\alpha$. The binormal vector is $B=T \wedge_{L} N$, which is a spacelike vector and the torsion of $\alpha$ is $\tau=\langle\dot{N}, B\rangle$.
The Bishop frame is an alternative approach to defining a moving frame that is well defined even when the spacelike curve with a spacelike binormal has vanishing second derivative. We can parallel transport an orthonormal frame along a spacelike curve with a spacelike binormal simply by parallel transporting each component of the frame. The parallel transport frame is based on the observation that, while $T(t)$ for a given spacelike curve with a spacelike binormal model is unique, we may choose any convenient arbitrary basis $\left(N_{1}(t), N_{2}(t)\right)$ for the remainder of the frame, so long as it is in the normal plane perpendicular to $T(t)$ at each point. If the derivatives of $\left(N_{1}(t), N_{2}(t)\right)$ depend only on $T(t)$ and not each other we can make $N_{1}(t)$ and $N_{2}(t)$ vary smoothly throughout the path regardless of the curvature. Therefore, we have the alternative frame equations:

$$
\begin{aligned}
& \dot{T}=\kappa_{1} N_{1}-\kappa_{2} N_{2}, \dot{N}_{1}=\kappa_{1} T, \dot{N}_{2}=\kappa_{2} T \\
& \langle T, T\rangle_{L}=1,\langle N, N\rangle_{L}=-1,\left\langle N_{2}, N_{2}\right\rangle_{L}=1 \\
& \kappa(t)=\sqrt{\left|\kappa_{1}^{2}-\kappa_{2}^{2}\right|}, \tau(t)=\frac{d \theta(t)}{d t}, \quad \theta(t)=\arctan h\left(\frac{\kappa_{2}}{\kappa_{1}}\right)
\end{aligned}
$$

$[2,3]$. So that $\kappa_{1}$ and $\kappa_{2}$ effectively correspond to a Cartesian coordinate system for the polar coordinates $\kappa, \theta$ with $\theta=\int \tau(t) d t$. The orientation of the parallel transport frame includes the arbitrary choice of integration constant $\theta_{0}$, which disappears from $\tau$ (and hence from the Frenet frame) due to the differentiation.

## 2 Introduction

In one parameter motion of a body in Lorentz-Minkowski 3 -space is generated by the transformation

$$
f: E_{1}^{3} \longrightarrow E_{1}^{3}
$$

$$
\begin{equation*}
X \longrightarrow f(X)=Y=A X+C \tag{1}
\end{equation*}
$$

Where $A \in S O_{1}(3)$ and $X, Y, C$ are $3 \times 1$ real matrices and

$$
S O_{1}(3)=\left\{A \in R_{3}^{3} \mid \operatorname{det} A=1, \quad A^{t} \varepsilon A=\varepsilon, \quad \varepsilon=\left[\begin{array}{ccc}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & -1
\end{array}\right]\right\}
$$

$A, C$ are $C^{\infty}$ functions of a real parameter $t, X$ and $Y$ corresponding to the position vectors of the same point $X$, with respect to the orthonormal coordinate systems of the moving space $H$ and the fixed space $H_{0}$, respectively. At the initial time $t=t_{0}$ we consider the coordinate system of $H_{0}$ and $H$ are coincident $[2,6]$.
In the special Frenet and Bishop motions, $A$ matrix is $[T N B]$ and $\left[T N_{1} N_{2}\right.$ ] respectively.
In this paper, we first find a geometrical meaning for $\operatorname{det} \dot{A}$, $\operatorname{det} \ddot{A}$ and $\operatorname{det} \dddot{A}$. The $1^{\text {st }}$ order velocity of a fixed point $X$ is $\dot{Y}=\dot{A} X+\dot{C}$ and for the $2^{\text {nd }}$ and $3^{\text {rd }}$ order velocity of this point, give us $\ddot{Y}=\ddot{A} X+\ddot{C}$ and $\dddot{Y}=\dddot{A} X+\dddot{C}$ respectively.
$\dot{Y}$ is the sliding velocity and $\ddot{Y}$ and $\dddot{Y}$ are the $1^{\text {st }}$ and $2^{\text {nd }}$ sliding acceleration of the point $X$ respectively. We will show that existence of the $1^{\text {st }}$ and $2^{\text {nd }}$ acceleration pole points by the solution of the $\ddot{A} X+\ddot{C}=0$ and $\dddot{A} X+\dddot{C}=0$ systems. The solution of these systems depend on $\operatorname{det} \ddot{A}$ and $\operatorname{det} \dddot{A}$.

## 3 Acceleration Pole Points In Frenet Motion

Definition 3.1 The first derivation of (1), with respect to $t$, we have

$$
\dot{Y}=\dot{A} X+\dot{C}+A \dot{X}
$$

Where $\dot{Y}$ is the absolute velocity, $\dot{A} X+\dot{C}$ is the sliding velocity and $A \dot{X}$ is the relative velocity of the point $X$. The solution vector $X$ of the system $\dot{A} X+\dot{C}=0$ is the position vector of the point which may be considered as a fixed point of $H_{0}$ and $H$ at the same time $t$. These points are called instantaneous pole points at the time $t$. The sliding velocity of a fixed point $X$ in moving space $H$ is

$$
\begin{equation*}
\dot{Y}=\dot{A} X+\dot{C} \tag{2}
\end{equation*}
$$

and for the $2^{n d}$ order velocity (or the $1^{\text {st }}$ order sliding acceleration) of this point, (2) gives us

$$
\begin{equation*}
\ddot{Y}=\ddot{A} X+\ddot{C} \tag{3}
\end{equation*}
$$

and for the $3^{\text {rd }}$ order velocity (or the $2^{\text {nd }}$ sliding acceleration) of this point, (3) gives us

$$
\begin{equation*}
\dddot{Y}=\dddot{A} X+\dddot{C} \tag{4}
\end{equation*}
$$

By using the Frenet formulas and

$$
A=\left[\begin{array}{lll}
T & N & B
\end{array}\right], \dot{A}=\left[\begin{array}{lll}
\dot{T} & \dot{N} & \dot{B}
\end{array}\right], \quad \ddot{A}=\left[\begin{array}{lll}
\ddot{T} & \ddot{N} & \ddot{B}
\end{array}\right], \quad \dddot{A}=\left[\begin{array}{lll}
\dddot{T} & \dddot{N} & \dddot{B}
\end{array}\right]
$$

we can give,

$$
\begin{aligned}
& {\left[\begin{array}{c}
\dot{T} \\
\dot{N} \\
\dot{B}
\end{array}\right]=\left[\begin{array}{lll}
0 & \kappa & 0 \\
\kappa & 0 & \tau \\
0 & \tau & 0
\end{array}\right]\left[\begin{array}{c}
T \\
N \\
B
\end{array}\right],} \\
& \operatorname{det} \dot{A}=\left|\begin{array}{ccc}
0 & \kappa & 0 \\
\kappa & 0 & \tau \\
0 & \tau & 0
\end{array}\right| \cdot \operatorname{det} A=0
\end{aligned}
$$

Then the system $\dot{A} X+\dot{C}=0$ has not unique solution. So, the Frenet motion has not pole point.

## $3.1 \quad 1^{\text {st }}$ acceleration pole points in Frenet motion

The discussion of existence of the $1^{\text {st }}$ acceleration poles and the $1^{\text {st }}$ acceleration axodes is the discussion of the solution of the system

$$
\begin{equation*}
\ddot{A} X+\ddot{C}=0 \tag{5}
\end{equation*}
$$

The solution of the system of (5) depends on $\operatorname{det} \ddot{A}$.
Theorem 3.2 The Spacelike curve with a spacelike binormal $\alpha(t)$ is not general helix; iff the Frenet motion has a $1^{\text {st }}$ acceleration pole point in the moving space $H ; X=-(\ddot{A})^{-1} \ddot{C}$.

Proof. If $\{T, N, B\}$ is an adapted Frenet frame, then we have

$$
\begin{aligned}
& \ddot{T}=\kappa^{2} T+\dot{\kappa} N+\kappa \tau B \\
& \ddot{N}=\dot{\kappa} T+\left(\kappa^{2}+\tau^{2}\right) N+\dot{\tau} N \\
& \ddot{B}=\kappa \tau T+\dot{\tau} N+\tau^{2} B
\end{aligned}
$$

So, we obtain

$$
\left[\begin{array}{c}
\ddot{T} \\
\ddot{N} \\
\ddot{B}
\end{array}\right]=\left[\begin{array}{ccc}
\kappa^{2} & \dot{\kappa} & \kappa \tau \\
\dot{\kappa} & \left(\kappa^{2}+\tau^{2}\right) & \dot{\tau} \\
\kappa \tau & \dot{\tau} & \tau^{2}
\end{array}\right]\left[\begin{array}{c}
T \\
N \\
B
\end{array}\right] .
$$

By using $A=[T N B] \in S O_{1}(3)$, we get

$$
\operatorname{det} \ddot{A}=\left|\begin{array}{ccc}
\kappa^{2} & \dot{\kappa} & \kappa \tau  \tag{6}\\
\dot{\kappa} & \left(\kappa^{2}+\tau^{2}\right) & \dot{\tau} \\
\kappa \tau & \dot{\tau} & \tau^{2}
\end{array}\right|=-\left[\kappa^{2}\left(\frac{\tau}{\kappa}\right)\right]^{2}
$$

Obviously as a consequence of equation (6) we have the following:

$$
\operatorname{det} \ddot{A}=0 \Leftrightarrow \frac{\tau}{\kappa}=\text { constant }
$$

From this case we obtain that at any moment $t$, if the curve $\alpha(t)$ is a generalized helix then the solution systems of (5) are not unique in fixed space $H_{0}$. The Frenet motion $Y=A X+C$ has not the $1^{\text {st }}$ acceleration pole point. If $\operatorname{det} \ddot{A} \neq 0$ then $\alpha(t)$ is not general helix and Frenet motion has a $1^{\text {st }}$ acceleration pole point, $X=-(\ddot{A})^{-1} \ddot{C}$.

## $3.22^{\text {nd }}$ acceleration pole points in Frenet motion

The discussion of existence of the $2^{\text {nd }}$ acceleration pole points and the $2^{\text {nd }}$ acceleration axodes is the discussion of the solution of the system

$$
\begin{equation*}
\dddot{A} X+\dddot{C}=0 \tag{7}
\end{equation*}
$$

Theorem 3.3 If the spacelike curve with a spacelike binormal curve $\alpha(t)$ is a generalized helix; then the Frenet motion has not a $2^{n d}$ acceleration pole point in fixed space $H_{0}$.

Proof. If $T, N$ and $B$ is an adapted Frenet frame, then we have;

$$
\begin{aligned}
& \dddot{T}=(3 \kappa \dot{\kappa}) T+\left(\kappa^{3}+\kappa \tau^{2}+\ddot{\kappa}\right) N+(\kappa \dot{\tau}+2 \dot{\kappa} \tau) B \\
& \dddot{N}=\left(\kappa^{3}+\kappa \tau^{2}+\ddot{\kappa}\right) T+3(\kappa \dot{\kappa}+\tau \dot{\tau}) N+\left(\tau^{3}+\kappa^{2} \tau+\ddot{\tau}\right) B \\
& \dddot{B}=(2 \kappa \dot{\tau}+\dot{\kappa} \tau) T+\left(\tau^{3}+\kappa^{2} \tau+\ddot{\tau}\right) N+(3 \tau \dot{\tau}) B
\end{aligned}
$$

So, we obtain

$$
\left[\begin{array}{c}
\dddot{T} \\
\dddot{N} \\
\dddot{B}
\end{array}\right]=\left[\begin{array}{ccc}
(3 \kappa \dot{\kappa}) & \left(\kappa^{3}+\kappa \tau^{2}+\ddot{\kappa}\right) & (\kappa \dot{\tau}+2 \dot{\kappa} \tau) \\
\left(\kappa^{3}+\kappa \tau^{2}+\ddot{\kappa}\right) & 3(\kappa \dot{\kappa}+\tau \dot{\tau}) & \left(\tau^{3}+\kappa^{2} \tau+\ddot{\tau}\right) \\
(2 \kappa \dot{\tau}+\dot{\kappa} \tau) & \left(\tau^{3}+\kappa^{2} \tau+\ddot{\tau}\right) & (3 \tau \dot{\tau})
\end{array}\right]\left[\begin{array}{l}
T \\
N \\
B
\end{array}\right] .
$$

By using $\operatorname{det} A=1$, we get

$$
\operatorname{det} \dddot{A}=3 \kappa^{2}\left(\frac{\tau}{\kappa}\right)^{\dot{2}}\left[-2 \kappa^{2}\left(\frac{\tau}{\kappa}\right)(\kappa \dot{\kappa}+\tau \dot{\tau})(\kappa \dot{\tau}-\dot{\kappa} \tau)+\left(\kappa^{2}+\tau^{2}\right)(\kappa \ddot{\tau}-\ddot{\kappa} \tau)\right]
$$

$$
\begin{equation*}
-3 \dot{\kappa}^{2}\left(\frac{\dot{\tau}}{\dot{\kappa}}\right)(\kappa \ddot{\tau}-\ddot{\kappa} \tau) \tag{8}
\end{equation*}
$$

As a consequence of equation of (7) we have the following:
Because $\alpha(t)$ is a general helix, then we can write

$$
\left(\frac{\tau}{\kappa}\right)=0 \text { and }\left(\frac{\dot{\tau}}{\dot{\kappa}}\right)=0
$$

Thus $\operatorname{det} \dddot{A}$.
From this case we obtain, at any time $t$, the curve $\alpha(t)$ is a generalized helix, and the solution of system (7) are not unique and in fixed space $H_{0}$, the Frenet motion $Y=A X+C$ has not the $2^{\text {nd }}$ acceleration pole point.

## 4 Acceleration Pole Points In Bishop Motion

By using the Bishop formulas and

$$
\begin{array}{lll}
A=\left[\begin{array}{lll}
T & N_{1} & N_{2}
\end{array}\right], & \dot{A}=\left[\begin{array}{lll}
\dot{T} & \dot{N}_{1} & \dot{N}_{2}
\end{array}\right], \\
\ddot{A}=\left[\begin{array}{lll}
\ddot{T} & \ddot{N}_{1} & \ddot{N}_{2}
\end{array}\right], & \dddot{A}=\left[\begin{array}{ccc}
\dddot{T} & N_{1} & N_{2}
\end{array}\right]
\end{array}
$$

we can give,

$$
\begin{aligned}
& {\left[\begin{array}{c}
\dot{T} \\
\dot{N}_{1} \\
\dot{N}_{2}
\end{array}\right]=\left[\begin{array}{ccc}
0 & \kappa_{1} & -\kappa_{2} \\
\kappa_{1} & 0 & 0 \\
\kappa_{2} & 0 & 0
\end{array}\right]\left[\begin{array}{c}
T \\
N_{1} \\
N_{2}
\end{array}\right],} \\
& \operatorname{det} \dot{A}=\left|\begin{array}{ccc}
0 & \kappa_{1} & -\kappa_{2} \\
\kappa_{1} & 0 & 0 \\
\kappa_{2} & 0 & 0
\end{array}\right|=0
\end{aligned}
$$

Then the system $\dot{A} X+\dot{C}=0$ has not unique solution. So, the Bishop motion has not pole point.

## $4.11^{\text {st }}$ acceleration pole points in Bishop motion

Theorem 4.1 The spacelike curve with a spacelike binormal $\alpha(t)$ is not planar in the moving space $H$ iff The Bishop motion has a $1^{\text {st }}$ acceleration pole point; $X=-(\ddot{A})^{-1} \ddot{C}$,

Proof. If $\left\{T, N_{1}, N_{2}\right\}$ is an adapted Bishop frame, then we have

$$
\begin{aligned}
& \ddot{T}=\left(\kappa_{1}^{2}-\kappa_{2}^{2}\right) T+\dot{\kappa}_{1} N_{1}-\dot{\kappa}_{2} N_{2} \\
& \ddot{N}_{1}=\dot{\kappa}_{1} T+\kappa_{1}^{2} N_{1}-\kappa_{1} \kappa_{2} N_{2} \\
& \ddot{N}_{2}=\dot{\kappa}_{2} T+\kappa_{1} \kappa_{2} N_{1}-\kappa_{2}^{2} N_{2}
\end{aligned}
$$

So, we obtain

$$
\left[\begin{array}{c}
\ddot{T} \\
\ddot{N}_{1} \\
\ddot{N}_{2}
\end{array}\right]=\left[\begin{array}{ccc}
\left(\kappa_{1}^{2}-\kappa_{2}^{2}\right) & \dot{\kappa}_{1} & -\dot{\kappa}_{2} \\
\dot{\kappa}_{1} & \kappa_{1}^{2} & -\kappa_{1} \kappa_{2} \\
\dot{\kappa}_{2} & \kappa_{1} \kappa_{2} & -\kappa_{2}^{2}
\end{array}\right]\left[\begin{array}{c}
T \\
N_{1} \\
N_{2}
\end{array}\right]
$$

By using $A=\left[\begin{array}{lll}T & N_{1} & N_{2}\end{array}\right]$, we get

$$
\operatorname{det} \ddot{A}=\left|\begin{array}{ccc}
\left(\kappa_{1}^{2}-\kappa_{2}^{2}\right) & \dot{\kappa}_{1} & -\dot{\kappa}_{2}  \tag{9}\\
\dot{\kappa}_{1} & \kappa_{1}^{2} & -\kappa_{1} \kappa_{2} \\
\dot{\kappa_{2}} & \kappa_{1} \kappa_{2} & -\kappa_{2}^{2}
\end{array}\right|=\left[\kappa_{1}^{2}\left(\frac{\kappa_{2}}{\kappa_{1}}\right)\right]^{2}
$$

By using equations:

$$
\kappa(t)=\sqrt{\left|\kappa_{1}^{2}-\kappa_{2}^{2}\right|,}, \quad \tau(t)=\frac{d \theta(t)}{d t}, \quad \theta(t)=\arctan h\left(\frac{\kappa_{2}}{\kappa_{1}}\right)
$$

we have,

$$
\begin{align*}
\kappa^{2}=\left|\kappa_{1}^{2}-\kappa_{2}^{2}\right| & \\
\frac{\kappa_{2}}{\kappa_{1}}=\tan h \theta & \Rightarrow\left(\frac{\kappa_{2}}{\kappa_{1}}\right)=\left(1-\tan ^{2} h \theta\right) \frac{d \theta}{d t} \\
& \Rightarrow\left(\frac{\kappa_{2}}{\kappa_{1}}\right)=\left(1-\frac{\kappa_{2}^{2}}{\kappa_{1}^{2}}\right) \tau=\left(\frac{\kappa_{1}^{2}-\kappa_{2}^{2}}{\kappa_{1}^{2}}\right) \tau \\
& \Rightarrow\left(\frac{\kappa_{2}}{\kappa_{1}}\right)= \pm\left(\frac{\kappa^{2}}{\kappa_{1}^{2}}\right) \tau \\
& \Rightarrow \kappa_{1}^{2}\left(\frac{\kappa_{2}}{\kappa_{1}}\right)= \pm \kappa^{2} \tau \tag{10}
\end{align*}
$$

Obviously as a consequence of equations (9) and (10) we have the following:

$$
\begin{equation*}
\operatorname{det} \ddot{A}=\kappa^{4} \tau^{2} \tag{11}
\end{equation*}
$$

As a consequence of equation of (11) we have the following:

$$
\operatorname{det} \ddot{A}=0 \Leftrightarrow \tau=0
$$

From this case we obtain, the solution systems of (5) are not unique in fixed space $H_{0}$ if and only if, at any time $t$, the curve $\alpha(t)$ is a plane. So that, the Bishop motion $Y=A X+C$ has not the $1^{\text {st }}$ acceleration pole point.
If $\operatorname{det} \ddot{A} \neq 0$ then $\alpha(t)$ is not plane.

## $4.22^{\text {nd }}$ acceleration pole points in Bishop motion

Theorem 4.2 The spacelike curve with a spacelike binormal $\alpha(t)$ is a plane $\Rightarrow$ in fixed space $H_{0}$, the Bishop motion has not a $2^{\text {nd }}$ acceleration pole point.

Proof. If $T, N_{1}$ and $N_{2}$ is an adapted Bishop frame, then we have;

$$
\begin{align*}
& \dddot{T}=3\left(\kappa_{1} \dot{\kappa}_{1}-\kappa_{2} \dot{\kappa}_{2}\right) T+\left(\kappa_{1}^{3}-\kappa_{1} \kappa_{2}^{2}+\ddot{\kappa}_{1}\right) N_{1}+\left(\kappa_{2}^{3}-\kappa_{1}^{2} \kappa_{2}-\ddot{\kappa}_{2}\right) N_{2} \\
& \ldots N_{1}=\left(\kappa_{1}^{3}-\kappa_{1} \kappa_{2}^{2}+\ddot{\kappa}_{1}\right) T+\left(3 \kappa_{1} \dot{\kappa}_{1}\right) N_{1}-\left(\kappa_{1} \dot{\kappa}_{2}+2 \dot{\kappa}_{1} \kappa_{2}\right) N_{2} \\
& \dddot{N_{2}}=\left(-\kappa_{2}^{3}+\kappa_{1}^{2} \kappa_{2}+\ddot{\kappa}_{2}\right) T+\left(\kappa_{2} \dot{\kappa}_{1}+2 \dot{\kappa}_{2} \kappa_{1}\right) N_{1}-\left(3 \kappa_{2} \dot{\kappa}_{2}\right) N_{2} \\
& {\left[\begin{array}{c}
\dddot{T} \\
\dddot{N} \\
N_{1} \\
N_{2}
\end{array}\right]=\left[\begin{array}{cc}
3\left(\kappa_{1} \dot{\kappa}_{1}-\kappa_{2} \dot{\kappa}_{2}\right) & \left(\kappa_{1}^{3}-\kappa_{1} \kappa_{2}^{2}+\ddot{\kappa}_{1}\right) \\
\left(\kappa_{2}^{3}-\kappa_{1}^{2} \kappa_{2}-\ddot{\kappa}_{2}\right) \\
\left(-\kappa_{2}^{3}-\kappa_{1} \kappa_{2}^{2}+\kappa_{1}^{2} \kappa_{2}+\ddot{\kappa}_{2}\right) & \left(\kappa_{2} \dot{\kappa}_{1}+2 \kappa_{1} \dot{\kappa}_{2} \kappa_{1}\right) \\
-\left(\kappa_{1} \dot{\kappa}_{2}+2 \dot{\kappa}_{1} \kappa_{2}\right) \\
-\left(3 \kappa_{2} \dot{\kappa}_{2}\right)
\end{array}\right]\left[\begin{array}{c}
T \\
N_{1} \\
N_{2}
\end{array}\right]} \\
& \qquad \begin{array}{r}
\operatorname{det} \dddot{A}=\operatorname{det}(\dddot{T}, \dddot{N}, \dddot{B}) \\
=3\left(2 \kappa \dot{\kappa} \tau+\kappa^{2} \dot{\tau}\right)\left(-\kappa^{4} \tau+\dot{\kappa}_{1}^{2}\left(\frac{\dot{\kappa}_{2}}{\dot{\kappa}_{1}}\right)\right)+6 \kappa^{5} \dot{\kappa} \tau^{2}
\end{array}
\end{align*}
$$

As a consequence of equation of (12) we have the following:

$$
\tau=0 \Rightarrow \operatorname{det} \dddot{A}=0
$$

From this case we obtain, if at any time $t$, the curve $\alpha(t)$ is a plane, then the solution of system (7) are not unique in fixed space $H_{0}$ and the Bishop motion $Y=A X+C$ has not the $2^{\text {nd }}$ acceleration pole point.

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